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COMBINING FAST MULTIPOLE TECHNIQUES AND AN APPROXIMATE INVERSE PRECONDITIONER FOR LARGE ELECTROMAGNETISM CALCULATIONS*

B. CARPENTIERI[†], I. S. DUFF[‡], L. GIRAUD[§], AND G. SYLVAND[¶]

Abstract. The boundary element method has become a popular tool for the solution of Maxwell's equations in electromagnetism. From a linear algebra point of view, this leads to the solution of large dense complex linear systems, where the unknowns are associated with the edges of the mesh defined on the surface of the illuminated object. In this paper, we address the iterative solution of these linear systems via preconditioned Krylov solvers. Our primary focus is on the design of an efficient parallelizable preconditioner. In that respect, we consider an approximate inverse method based on the Frobenius-norm minimization. The preconditioner is constructed from a sparse approximation of the dense coefficient matrix, and the patterns both for the preconditioner and for the coefficient matrix are computed a priori using geometric information from the mesh. We describe how such a preconditioner can be naturally implemented in a parallel code that implements the multipole technique for the matrix-vector product calculation. We investigate the numerical scalability of our preconditioner on realistic industrial test problems and show that it exhibits some limitations on very large problems of size close to one million unknowns. To improve its robustness on those large problems we propose an embedded iterative scheme that combines nested GMRES solvers with different fast multipole computations. We show through extensive numerical experiments that this new scheme is extremely robust at affordable memory and CPU costs for the solution of very large and challenging problems.

Key words. sparse approximate inverse preconditioner, large dense complex systems, Maxwell equations, boundary elements, fast multipole technique, frequency domain, electric field integral equations, parallel computing, large calculations

AMS subject classifications. 65F10, 65R20, 65N38, 78A55, 78M15

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1. Introduction. The analysis of wave propagation phenomena is gaining an increasing interest in recent years in the simulation of many challenging industrial processes, including the prediction of the Radar Cross Section (RCS) of arbitrarily shaped 3D objects such as aircraft, the study of electromagnetic compatibility of electrical devices with their environment, the design of antennae and absorbing materials, and many others. All these simulations are very demanding in terms of computer resources and require fast and efficient numerical methods to compute an approximate solution of Maxwell's equations. Using the equivalence principle, Maxwell's equations

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can be recast in the form of four integral equations that relate the electric and magnetic fields to the equivalent electric and magnetic currents on the surface of the object.

Among integral formulations, the electric field integral equation (EFIE) is the most general for electromagnetic scattering problems, as it can handle fairly general geometries, and thus is widely used in industrial simulations. The EFIE provides a first-kind integral equation, which is well known to be ill-conditioned and gives rise to linear systems that are challenging to solve by iterative methods. The discretization is performed on the surface of the object and gives rise to a linear system

$$(1.1) \quad Ax = b,$$

where the matrix A is an $n \times n$, dense, complex, symmetric, non-Hermitian matrix. For many years, direct solution methods have been the methods of choice for solving systems (1.1) because they are reliable and predictable in terms of both accuracy and computing costs. However, for the solution of large-scale problems, direct methods are infeasible even on large parallel platforms because they require the unaffordable storage of n^2 single or double precision complex entries of the coefficient matrix and $\mathcal{O}(n^3)$ floating-point operations to compute the factorization. The use of preconditioned Krylov solvers can be an alternative to direct solution methods, provided we have fast matrix-free matrix-vector products and robust preconditioners. Research efforts have recently concentrated on fast methods for performing matrix-vector products with $\mathcal{O}(n \log(n))$ computational complexity, including strategies for parallel distributed memory implementations. These methods, generally referred to as *hierarchical methods*, were introduced originally in the context of the study of particle simulations and can be used effectively in boundary element applications.

In this paper, we focus on the design of a parallelizable preconditioner to be used in conjunction with a parallel distributed fast multipole technique that implements the matrix-vector calculation. We consider an approximate inverse preconditioner based on Frobenius-norm minimization with a pattern prescribed in advance. In section 2, we describe the implementation of the preconditioner within an out-of-core parallel code that implements the fast multipole method (FMM) for the matrix-vector product. We investigate the numerical scalability of the preconditioner on a set of industrial problems of increasing size and show that it becomes less effective when the problem size becomes very large. To overcome this weakness, we propose in section 3 an embedded iterative scheme based on the GMRES method that aims at improving the robustness of the preconditioner on large applications. We illustrate the numerical efficiency and the cost effectiveness of the proposed scheme on systems of up to one million unknowns arising from challenging problems from the electromagnetism community.

2. Preconditioning boundary integral equations. For large meshes with many surface details, the density of the discretization mesh is nonuniform, and the matrix generated by the method of moments can become ill-conditioned. The convergence of Krylov methods depends to a large extent on the eigenvalue distribution of the coefficient matrix. In Figure 2.1, we plot the eigenvalue distribution in the complex plane of the matrix associated with a satellite geometry. This distribution is representative of the general trend; it can be seen that the eigenvalues of the system are very scattered; many of them have a large negative real part and no clustering appears. Such a distribution is not at all favorable for the rapid convergence of Krylov methods. One goal of preconditioning is to improve this distribution by grouping the

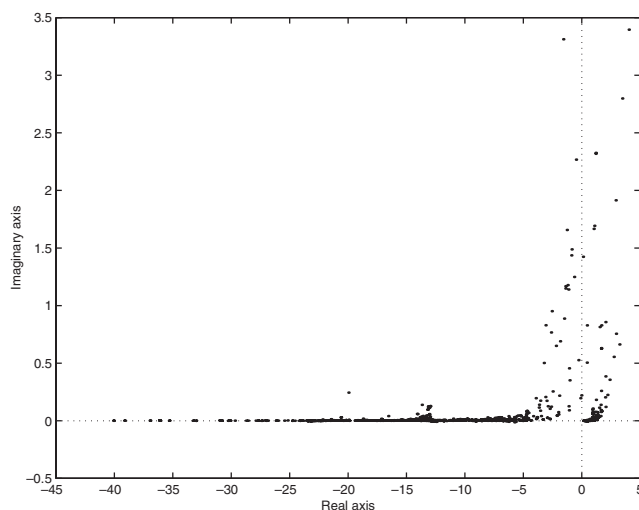


FIG. 2.1. *Eigenvalue distribution in the complex plane of the coefficient matrix for a scattering problem from a satellite that is representative of the general trend.*

eigenvalues into a few small clusters.

The design of robust preconditioners for boundary integral equations can be challenging. Simple preconditioners like the diagonal of A , diagonal blocks, or a band can be effective only when the coefficient matrix has some degree of diagonal dominance arising from the integral formulation [30]. Block diagonal preconditioners are generally more robust than their pointwise counterparts but may require matrix permutations or renumbering of the grid points to cluster the large entries close to the diagonal. Incomplete factorizations have been successfully used on nonsymmetric dense systems in [28] and hybrid integral formulations in [24], but for the EFIE the triangular factors computed by the factorization are often very ill-conditioned due to the indefiniteness of A . This makes the triangular solves highly unstable and the preconditioner useless. In [5], we used a small diagonal shift applied to A before computing the factorization with the intention of moving the eigenvalues along the imaginary axis. This might avoid a possible eigenvalue cluster close to zero and can help the computation of more stable factors in some cases. However, this shift is not easy to tune and there is no way to predict its effect.

Approximate inverse methods are generally less prone to instabilities on indefinite systems, and several preconditioners of this type have been proposed in electromagnetism (see for instance [1, 3, 4, 6, 23, 27, 35]). Owing to the rapid decay of the discrete Green's function, the location of the large entries in the inverse matrix exhibits some structure. In addition, only a few of its entries have relatively large magnitude. This means that a very sparse matrix is likely to retain the most relevant contributions to the exact inverse. This remarkable property can be effectively exploited in the design of robust approximate inverses as preconditioners for electromagnetism applications.

2.1. Frobenius-norm minimization preconditioner. In this section, we describe an approximate inverse preconditioner based on Frobenius-norm minimization. The original idea, due to Benson and Frederickson [2, 15], is to compute the sparse approximate inverse as the matrix M which minimizes $\|I - MA\|_F$ (or $\|I - AM\|_F$ for right preconditioning) subject to certain sparsity constraints. The Frobenius norm

is usually chosen since it allows for the decoupling of the minimization problems into n independent linear least-squares problems, one for each column of M , when preconditioning from the right (or row of M , when preconditioning from the left). The independence of these least-squares problems follows immediately from the identity

$$(2.1) \quad \|I - AM\|_F^2 = \sum_{j=1}^n \|e_j - Am_{\bullet j}\|_2^2,$$

where e_j is the j th canonical unit vector, $m_{\bullet j}$ is the column vector representing the j th column of M , and n the dimension of the square matrices. In the case of left preconditioning, because A is symmetric, the analogous relation

$$(2.2) \quad \|I - MA\|_F^2 = \|I - AM^T\|_F^2 = \sum_{j=1}^n \|e_j - Am_{j\bullet}\|_2^2$$

holds, where $m_{j\bullet}$ is the column vector representing the j th row of M . Clearly, there is considerable scope for parallelism in this approach. The main issue is the selection of the sparsity pattern of M , that is, the set of indices

$$(2.3) \quad S = \{(i, j) \subseteq [1, n]^2 \text{ such that } m_{ij} = 0\}.$$

The idea is to keep M reasonably sparse while trying to capture the “large” entries of the inverse, which are expected to contribute the most to the quality of the preconditioner. Two different approaches can be followed for this purpose: an adaptive technique that dynamically tries to identify the best structure for M , and a static technique, where the pattern of M is prescribed a priori based on some heuristics. Adaptive methods usually start with a simple initial guess, like a diagonal matrix, and then improve the pattern until a criterion of the form $\|Am_{\bullet j} - e_j\|_2 < \varepsilon$ (for each j) is satisfied for a given $\varepsilon > 0$ or until a maximum number of nonzeros in $m_{\bullet j}$ is reached. If the norm is larger than ε and the number of nonzeros used is less than a fixed maximum, the pattern is enlarged according to some heuristics and the j th column of the approximate inverse is recomputed. The process is repeated until the required accuracy or storage limit is met (see [7, 21]).

Adaptive strategies can solve fairly general and hard problems but tend to be very expensive. The use of effective static pattern selection strategies can greatly reduce the amount of work in terms of CPU time and can substantially improve the overall setup process. When the coefficient matrix has a special structure or special properties, efforts have been made to find a pattern that can retain the entries of A^{-1} having large modulus [10, 11, 12, 32]. If A is row diagonally dominant, then the entries in the inverse decay columnwise and vice versa [32]. On boundary integral equations the discrete Green’s function decays rapidly far from the diagonal, and the inverse of A may have a very similar structure to that of A . The discrete Green’s function can be considered as a row or as a column of the exact inverse depicted on the physical computational grid. In this case a good pattern for the preconditioner can be computed in advance using graph information from \tilde{A} , a sparse approximation of the coefficient matrix constructed by dropping all the entries lower than a prescribed global threshold [1, 4, 22]. When fast methods are used for the matrix-vector products, all the entries of A are not available and the pattern can be formed by exploiting the near-field part of the matrix that is explicitly computed and available in the FMM [23].

Since we work in an integral equation context, relevant information for the construction of the pattern of M can be extracted from the mesh. When the object

geometries are smooth, only the neighboring edges (in the mesh topology sense) in the mesh can have a strong interaction with each other, while faraway connections are generally much weaker. Thus an effective pattern for the j th column of the approximate inverse can be computed by selecting in the mesh edge j and its q th level nearest neighbors. Three levels generally provide a good pattern for constructing an effective sparse approximate inverse. Using more levels increases the computational cost but does not improve the quality of the preconditioner substantially [4]. When the object geometries are not smooth or have disconnected parts, faraway edges in the mesh can have a strong interaction and be strongly coupled in the inverse matrix. In this case, a more robust pattern for the preconditioner can be computed using geometrical information, that is, selecting for each edge all those edges within a sufficiently large geometric neighborhood. In [4] we compared pattern selection strategies based on both algebraic and mesh information on a large set of problems and found that those exploiting geometric information are the most effective in capturing the large entries of the inverse.

In order to preserve sparsity, $\mathcal{O}(1)$ nonzero locations are computed in the pattern of the column $m_{\bullet j}$ of M . This makes each QR least-squares solution for each column of M from (2.1) cost $\mathcal{O}(n)$, and the overall construction of M costs approximately $\mathcal{O}(n^2)$ arithmetic operations. This cost can be significantly reduced if the preconditioner is computed using as input a sparse approximation \tilde{A} of the dense coefficient matrix A . On general problems, this approach can cause a severe deterioration of the quality of the preconditioner. In an integral equation context, it is likely to be more effective because the boundary element method generally introduces a very localized strong coupling among the edges in the underlying mesh. It means that a very sparse matrix can still retain the most relevant contributions from the singular integrals that give rise to dense matrices. If the sparsity pattern S of M is known in advance, the nonzero structure for the j th column of M is automatically determined and defined as

$$J = \{i \in [1, n] \text{ such that } (i, j) \in S\}.$$

The least-squares solution involves only the columns of \tilde{A} indexed by J ; we indicate this subset by $\tilde{A}(:, J)$. When \tilde{A} is sparse, many rows in $\tilde{A}(:, J)$ are usually null, not affecting the solution of the least-squares problems (2.1). Thus if I is the set of indices corresponding to the nonzero rows in $\tilde{A}(:, J)$, and if we define $\hat{A} = \tilde{A}(I, J)$, $\hat{m}_j = m_j(J)$, and $\hat{e}_j = e_j(J)$, the actual “reduced” least-squares problems are

$$(2.4) \quad \min \|\hat{e}_j - \hat{A}\hat{m}_j\|_2, \quad j = 1, \dots, n.$$

Usually problems (2.4) have much smaller size than problems (2.1) and can be efficiently solved by a dense QR factorization. In [1] the same nonzero sparsity pattern is selected for both \tilde{A} and M ; in that case, especially when the pattern is very sparse, the computed preconditioner may be poor on some geometries. Selecting more entries in \tilde{A} than in M can provide a more robust preconditioner, and the additional cost in terms of CPU time is negligible because of the complexity of the QR factorization [4]. Increasing the number of rows q_I , that is, the number of entries of \tilde{A} , affects the CPU time far less than increasing the density of the preconditioner, that is, the number of columns p_J in the least-squares problems. This is because each least-squares solution costs $\mathcal{O}(q_I p_J^2)$.

2.2. Implementation of the preconditioner in the FMM context. The FMM, introduced by Greengard and Rokhlin in [20], provides an algorithm for computing approximate matrix-vector products for electromagnetic scattering problems.

The method is fast in the sense that the computation of one matrix-vector product costs $\mathcal{O}(n \log n)$ arithmetic operations instead of the usual $\mathcal{O}(n^2)$ operations. It requires storage only of the near-field part of the matrix, that is, of $\mathcal{O}(n \log n)$ instead of $\mathcal{O}(n^2)$ entries. These properties mean that using iterative solvers for the solution of large problems becomes feasible.

The basic idea of the algorithm is to compute interactions among degrees of freedom in the mesh at different levels of accuracy depending on their physical distance. Single and multilevel variants of the FMM exist. The 3D object is first entirely enclosed in a large cube that is subdivided into eight cubes. Each cube is recursively divided until the size is small compared with the wavelength. In the hierarchical multilevel algorithm, the box-wise partitioning of the obstacle is carried out until the size of the smallest box is generally half of a wavelength, and tree-structured data is used at all levels. In particular, only nonempty cubes are indexed and recorded in the data structure. The resulting tree is called the *oct-tree* (see Figure 2.2) and its leaves are referred to as *leaf boxes*. The oct-tree provides a hierarchical representation of the computational domain partitioned by boxes: each cube has up to eight children and one parent in the oct-tree, except for the largest cube which encloses the whole domain. Obviously, the leaf boxes have no children. Multipole coefficients are computed for all cubes starting from the lowest level of the oct-tree, that is, from the leaf boxes, and then recursively for each parent cube by summing together multipole coefficients of its children. For each observation cube, an interaction list is defined which consists of those cubes that are not neighbors of the cube itself but whose parent is a neighbor of the cube's parent (see Figure 2.3 for a 2D representation). The interactions of degrees of freedom within neighboring boxes are computed exactly, while the interactions between cubes that are in the interaction list are computed using the FMM. All the other interactions are computed hierarchically on a coarser level by traversing the oct-tree. Both the computational cost and the memory requirement of the algorithm are of order $\mathcal{O}(n \log n)$. Further information on the algorithmic steps and recent theoretical investigations of the FMM can be found in [8, 9, 31]; also see [17, 19, 36] for discussions on parallel implementation issues.

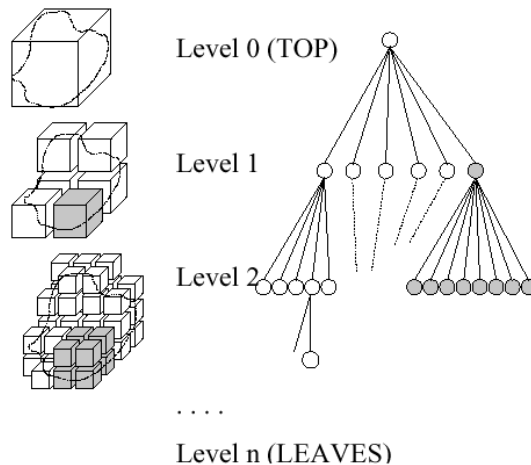


FIG. 2.2. The oct-tree in the FMM algorithm. The maximum number of children is eight. The actual number corresponds to the subset that intersects the object.

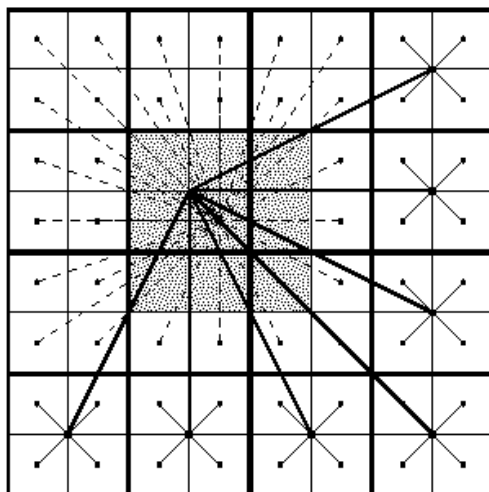


FIG. 2.3. *Interactions in the multilevel FMM. The interactions for the gray boxes are computed directly. We denote by dashed lines the interaction list for the observation box that consists of those cubes that are not neighbors of the cube itself but whose parent is a neighbor of the cube's parent. The interactions of the cubes in the list are computed using the FMM. All the other interactions are computed hierarchically on a coarser level, denoted by solid lines.*

The box-wise decomposition of the domain required by the FMM naturally leads to an a priori pattern selection strategy for M and \tilde{A} using geometric information, that is on the spatial distribution of its degrees of freedom. We will adopt the following criterion: the nonzero structure of the column of the preconditioner associated with a given edge is defined by retaining all the edges within its leaf box and those in one level of neighboring boxes, and the structure for the sparse approximation of the dense coefficient matrix is defined by retaining the entries associated with edges included in the given leaf box as well as those belonging to the two levels of neighbors. The approximate inverse has a sparse block structure; each block is dense and is associated with one leaf box. Indeed the least-squares problems corresponding to edges within the same box are identical because they are defined using the same nonzero structure and the same set of entries of A . It means that we need only compute one QR factorization per leaf box. In our implementation we use two different oct-trees, and thus two different partitionings, to assemble the approximate inverse and for the approximate multipole coefficient matrix. The size of the smallest boxes in the partitioning associated with the preconditioner is a user-defined parameter that can be tuned to control the number of nonzeros computed per column, that is, the density of the preconditioner. According to our criterion, the larger the size of the leaf boxes, the larger the geometric neighborhood that determines the sparsity structure of the columns of the preconditioner. Parallelism can be exploited by assigning disjoint subsets of leaf boxes to different processors and performing the least-squares solutions independently on each processor. We refer to [31] for a complete description of the parallel code that we use.

2.3. Numerical scalability of the preconditioner. In this section, we study the numerical scalability of the Frobenius-norm minimization preconditioner. The optimal behavior would be to get frequency independent numerical behavior that

would result in convergence independent of the problem size. We show results on the following geometries:

- an industrial civil aircraft from a European company (see Figure 2.4(a)), a real-life model problem in an industrial context;
- an almond (see Figure 2.4(b)) that is also a typical test case in electromagnetic simulations;
- a Cetaf (see Figure 2.4(c)), a typical test case in electromagnetic simulations;
- a Cobra (see Figure 2.4(d)), a complex geometry that represents an air intake.

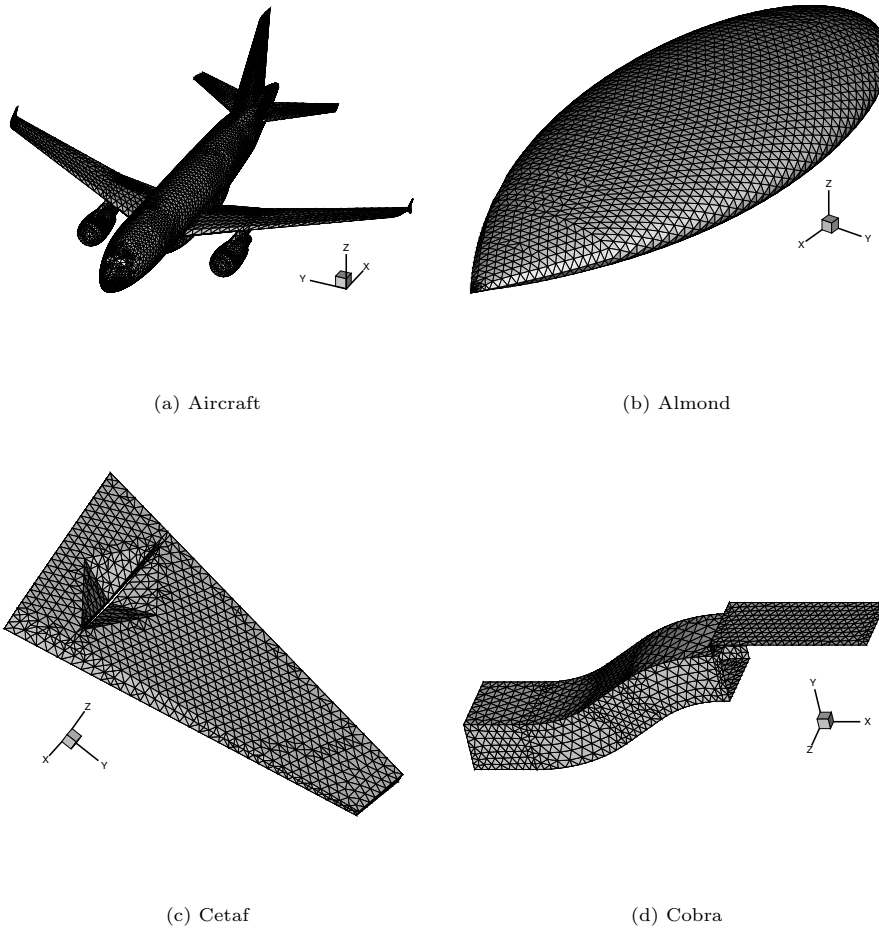


FIG. 2.4. *Mesh associated with test examples.*

In all the numerical experiments, the surface of the object is always discretized using 10 points per wavelength; larger discretizations are obtained by increasing the frequency of the illuminating wave. In all the experiments, we consider a right preconditioned GMRES method [26] and the threshold for the stopping criterion is set to 10^{-3} on the normwise backward error $\frac{\|r\|}{\|b\|}$, where r denotes the residual and b the right-hand side of the linear system. This tolerance is accurate for engineering

purposes, as it enables the correct construction of the radar cross section of the object. The initial guess is the zero vector. All the runs have been performed in single precision on eight processors of a Compaq Alpha server. The Compaq Alpha server is a cluster of symmetric multiprocessors. Each node consists of four DEC Alpha processors (EV 6, 1.3 GFlops peak) that share 512 MB of memory. On that computer, the temporary disk space that can be used by the out-of-core solver is around 189 GB. Among all the possible right-hand sides for each geometry, we have selected those that are the most difficult to solve in order to better illustrate the robustness and the efficiency of our preconditioner.

In Table 1, we show the numerical behavior of the preconditioner observed on the four geometries when the size of the problems is increased. In this table “d” means day, “h” hour, and “m” minute. We see that the numerical behavior of the preconditioner does not scale well with the size of the problems, especially GMRES(120) on the Cetaf, and there is no convergence on the largest aircraft problems. For GMRES(∞) the increase in the iteration count is less significant, even though on the Cetaf and the aircraft convergence cannot be obtained because we exceed either the memory limits of our computer or the time limit allocated to a single run. From a timing point of view, we see that the solution time of full GMRES is strongly affected by the orthogonalization involved in the Arnoldi procedure. On the Cetaf problem discretized with 531900 points, the number of iterations of GMRES(120) is twice as large with

TABLE 1

Number of matrix-vector products and elapsed time required to converge on the four problems on eight processors of the Compaq machine, except those marked with $^{(k)}$, that were run on k processors. Tolerance for the iterative solution was 10^{-3} . Acronyms: N.A. \equiv not available. M.L.E. \equiv memory limits exceeded.

Aircraft						
Size	Density	FROB	Time	FROB	GMRES(∞)	
					Iter	Time
94704	0.28		11m		746	2h 9m
213084	0.13		31m		973	7h 19m
591900	0.09		1h 30m		1461	16h 42m ⁽⁶⁴⁾
1160124	0.02		3h 24m		M.L.E. ⁽⁶⁴⁾	N.A.
					+2000	> 4d
Almond						
Size	Density	FROB	Time	FROB	GMRES(∞)	
					Iter	Time
104793	0.19		6m		234	20m
419172	0.05		21m		413	2h 44m
943137	0.02		49m		454	3h 35m ⁽³²⁾
					589	5h 55m
Cetaf						
Size	Density	FROB	Time	FROB	GMRES(∞)	
					Iter	Time
86256	0.18		4m		656	1h 25m
134775	0.11		6m		618	1h 45m
264156	0.06		13m		710	9h
531900	0.03		20m		844	1d 18m
1056636	0.01		37m		+750	+9h ⁽³²⁾
					+2000	> 1d
Cobra						
Size	Density	FROB	Time	FROB	GMRES(∞)	
					Iter	Time
60695	0.24		2m		369	26m
179460	0.09		7m		353	1h 11h
					406	1h 2m

respect to full GMRES, but GMRES(120) is about twice as cheap. Provided we get convergence, the use of a large restart often reduces the solution time even though it significantly deteriorates the convergence. On the Cetaf geometry, the solution time for the GMRES method increases superlinearly for small and medium problems, but nearly quadratically for large problems. On the largest test case, discretized with one million unknowns, un restarted GMRES does not converge after 750 iterations requiring more than nine hours of computation on 32 processors. The aircraft is very difficult to solve because the mesh has many surface details and the discretization matrices become ill-conditioned. On small and medium problems, the number of GMRES iterations increases with the problem size, and the solution time increases superlinearly. On the largest test case, discretized with one million unknowns, full GMRES exceeds the memory limit on 64 processors. In this case, the use of large restarts (120 in this table) does not enable convergence within 2000 iterations except on a small mesh of size 94704.

We see in Table 1 that our strategy for adjusting the leaf box dimension for increasing frequencies causes the density of the sparse approximate inverse to decrease for increasing problem size. The number of unknowns per box remains constant, but the number of boxes increases, leading to a decrease in density and to poorer preconditioners. In Table 2, we investigate the influence of the density on the quality of the preconditioner on the Cobra using GMRES with a large restart (120). We adopt the same criterion described in section 2 to define the sparsity patterns, but we increase the size of the leaf boxes in the oct-tree associated with the preconditioner. The best trade-off between cost and performance is obtained for a radius of around 0.12 wavelengths, the default value set in the code. If the preconditioner is used to solve systems with the same coefficient matrix and multiple right-hand sides, it might be worth computing more nonzeros if we have sufficient disk space, because the construction cost can be quickly amortized. However, significantly enlarging the density of the preconditioner is not feasible on the largest problems because we would exceed the memory and disk capacity of our computer. As mentioned already, our main focus is on linear systems arising from EFIE. The advantages of this formulation are numerous; in particular, it does not require any hypothesis on the geometry of the objects. We should nevertheless mention that for closed geometries the CFIE (combined field integral equation) can also be used. The linear systems arising from the CFIE formulation are much easier to solve. For instance the solution of the problem associated with the aircraft with 213084 degrees of freedom requires only 129 iterations of unpreconditioned full GMRES, and 22 iterations of preconditioned full GMRES. Furthermore, preconditioned full GMRES converges in 24 iterations on the aircraft with more than a million degrees of freedom. Because the linear systems arising from the CFIE are not challenging from a linear algebra point of view, we do not consider them further in this paper.

Finally, in Table 3, we show the parallel scalability of the implementation of the FMM code [31]. We solve problems of increasing size on a larger number of processors, keeping the number of unknowns per processor constant. It can be seen that the construction of the preconditioner scales perfectly. This is a benefit from keeping the size of the leaf box constant. Its use requires some communication but it still scales reasonably well. The scalability of the matrix-vector product is also satisfactory, as the increase in the elapsed time is due not only to the amount of data exchanges but also to the $\log(n)$ effect of its computational complexity.

TABLE 2

Number of matrix-vector products and elapsed time to build the preconditioner and to solve the problem using GMRES(120) on a Cobra problem of size 179460, varying the parameters controlling the density of the preconditioner. The runs were performed on eight processors of the Compaq machine.

Radius	Density	# mat-vec in GMRES(120)	Construction time	Solution time	Overall time
0.08	0.039	430	218	3562	3780
0.10	0.062	383	294	3008	3302
0.12	0.091	406	434	2954	3388
0.14	0.120	380	640	2873	3513
0.18	0.204	351	1863	3122	4985
0.24	0.358	280	4947	2583	7531

TABLE 3

Tests on the parallel scalability of the code with respect to the construction and application of the preconditioner and to the matrix-vector product operation on problems of increasing size. The test example is the aircraft.

Problem size	Nb procs	Construction time (sec)	Elapsed time precond (sec)	Elapsed time mat-vec (sec)
112908	8	513	0.39	1.77
161472	12	488	0.40	1.95
221952	16	497	0.43	2.15
288300	20	520	0.45	2.28
342732	24	523	0.47	3.10
393132	28	514	0.47	3.30
451632	32	509	0.48	2.80
674028	48	504	0.54	3.70
900912	64	514	0.60	3.80

3. Improving the preconditioner robustness using embedded iterations.

The numerical results shown in the previous section indicate that the Frobenius-norm minimization preconditioner tends to become less effective when the problem size increases, especially on difficult problems. By its nature the sparse approximate inverse is inherently local because each degree of freedom is coupled to only a very few neighbors. Because the exact inverse is dense the compact support used to define the preconditioner may not allow an exchange of global information and on large problems the lack of a global approximation may have a severe impact on the convergence. In our implementation, the overall number of computed nonzeros decreases for increasing values of the frequency. When the preconditioner becomes very sparse, information related to the far-field is completely lost. In this case, some suitable mechanism has to be introduced to recover global information for the numerical behavior of the discrete Green's function.

In this section, we describe an embedded iterative scheme, combined with multipole techniques, that is designed to meet the goals of robustness, scalability, and parallelism of the iterative solver. The basic idea is to carry out a few steps of an inner Krylov method for the preconditioning operation. The overall algorithm results in the inner-outer scheme depicted in Figure 3.1. The outer solver must be able to work with variable preconditioners. Amongst various possibilities, we mention FGMRES [25] and GMRES* [33, p. 91]; this latter reduces to GMRESR [34] when the inner solver is GMRES. The efficiency of the proposed algorithm relies on two main factors: the inner solver has to be preconditioned so that the residual in the inner iterations can be significantly reduced in a few steps, and the matrix-vector

products within the inner and the outer solvers can be carried out with a different accuracy. The motivation that naturally leads us to consider inner-outer schemes is to try to balance the locality of the preconditioner with the use of the multipole approach. Experiments conducted in [18] with inner-outer schemes combined with multipole techniques on the potential equation were not entirely successful. In that case, no preconditioner was used in the inner solver. The desirable feature of using different accuracies for the matrix-vector products is enabled by the use of the FMM. In our scheme, a highly accurate FMM is used within the outer solver, as it governs the final accuracy of the computed solution. A less accurate FMM is used within the inner solver, as it is a preconditioner for the outer scheme; in this respect the inner iterations attempt only to give a rough approximation of the solution and consequently do not require an accurate matrix-vector calculation. In fact, we solve a nearby system for the preconditioning operation that enables us to save considerable computational effort during the iterative process, as the less accurate matrix-vector calculation requires about half of the computing time of the accurate one. In Table 4, we show the average elapsed time in seconds observed on eight processors for a matrix-vector product using the FMM with different accuracy levels.

```

Outer solver  $\longrightarrow$  FGMRES, GMRES $\star$ 

Do k = 1, 2, ...
  • Matrix-vector product: FMM with high accuracy
  • Preconditioning: Inner solver (GMRES, TFQMR, ...)
    Do i = 1, 2, ...
      • Matrix-vector product: FMM with low accuracy
      • Preconditioning:  $M_{Frob}$ 
    End Do
End Do

```

FIG. 3.1. Inner-outer solution schemes in the FMM context. Sketch of the algorithm.

TABLE 4

Average elapsed time observed on eight processors in seconds for a matrix-vector product using the FMM with different levels of accuracy.

Geometry	Inner FMM	Outer FMM
Aircraft 94704	3.5	4.8
Aircraft 213084	6.9	11.8
Aircraft 591900	17.4	30.2
Aircraft 1160124	34.5	66.5
Almond 104793	1.8	3.0
Almond 419172	6.2	11.2
Almond 943137	14.3	25.5
Cetaf 86256	1.9	3.7
Cetaf 134775	3.2	5.2
Cetaf 264159	5.6	11.4
Cetaf 539100	11.5	19.3
Cetaf 1056636	20.6	38.2
Cobra 60695	1.1	2.0
Cobra 179460	3.2	5.6

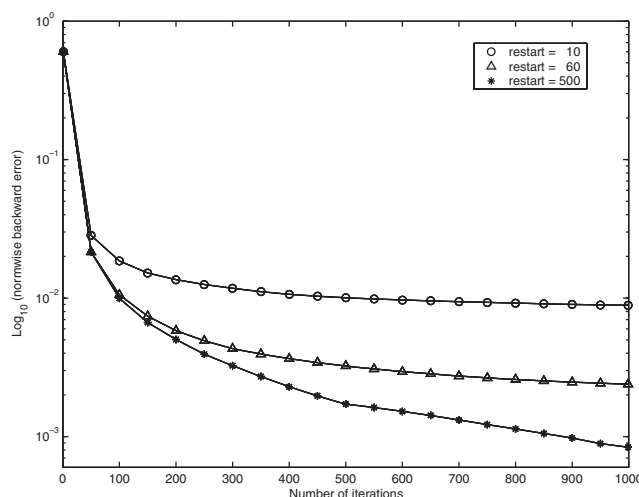


FIG. 3.2. Convergence history of GMRES for different values of restarts on an aircraft discretized with 94704 points.

3.1. Numerical results. In this section, we conduct experiments using the FGMRES method as the outer solver with an inner GMRES iteration preconditioned with the Frobenius-norm minimization method described in section 2. It is possible to vary the accuracy of the inner solver as the solution converges. However, in this paper, we present preconditioning results for a constant number of inner GMRES iterations. For the inner scheme we do not restart and perform only a prescribed number of full GMRES iterations or equivalently one step of restarted GMRES with the same prescribed restart. For the GMRES and FGMRES methods, we consider the implementations described in [14] and [13], respectively. The convergence history of GMRES depicted in Figure 3.2 for different values of the restart gives us some clues as to the numerical behavior of the proposed scheme. The residual of GMRES tends to decrease very rapidly in the first few iterations independently of the restarts, then decreases much more slowly, and finally tends to stagnate to a value that depends on the restart; the larger the restart, the lower the stagnation value. It suggests that a few steps in the inner solver could be very effective for obtaining a significant reduction of the initial residual. A different convergence behavior was observed when using other Krylov methods as an inner solver, for example, the TFQMR solver [3]. In this case, the residual at the beginning of the convergence is nearly constant or decreases very slowly. The use of TFQMR as an inner solver is ineffective.

In Table 5, we describe the results of experiments on an aircraft with 213084 degrees of freedom using different combinations of restarts for the inner and outer solvers. If the number of inner steps is too small, the preconditioning operation is poor and the convergence slows down, while too large restarts of GMRES tend to increase the overall computational cost but do not cause further reduction of the normwise backward error at the beginning of convergence. The choice of the restart for the outer solver depends to a large extent on the available memory of the target machine and the difficulty of the problem at hand, an issue that is also related to the illuminating direction of the incident wave that defines the right-hand side. Among the various possibilities, we select FGMRES(30) and GMRES(60) on the aircraft and

TABLE 5

Global elapsed time and total number of matrix-vector products required to converge on an aircraft with 213084 points varying the inner restart parameter. The runs have been performed on eight processors of the Compaq machine.

Restart FGMRES	Restart GMRES	# total inner mat-vec	# total outer mat-vec	Times
30	40	1960	51	4h 58m
30	50	1900	40	4h 53m
30	60	1920	34	4h 55m
30	70	2030	30	5h 13m
30	80	2240	29	5h 50m

the Cetaf problem, and FGMRES(15) and GMRES(30) on the Cobra and the Almond. These combinations are a good trade-off.

In Table 6 we show the results of experiments on the same geometries that we considered in section 2. We show the number of inner and outer matrix-vector products and the elapsed time needed to achieve convergence using a tolerance of 10^{-3} on eight processors of the Compaq machine. For comparison purposes, we show again in this table the results obtained with the restarted GMRES method. The comparison between the two solvers made in the tables is fair because GMRES has exactly the same storage requirements as the combination FGMRES/GMRES. In fact, for the same restart value, the storage requirements for the FGMRES algorithm are twice that for the standard GMRES algorithm, because it also stores the preconditioned vectors of the Krylov basis. It can be seen that the combination FGMRES/GMRES remarkably enhances the robustness of the preconditioner especially on large problems. The increase in the number of outer iterations is fairly modest except on the largest aircraft test cases. Nevertheless, the scheme is the only one that enables us to get the solution of this challenging problem since classical restarted GMRES does not converge and full GMRES exceeds the memory of our computer. Similarly, on the Cetaf discretized with one million points, the embedded scheme enables us to get convergence in 22 outer iterations, whereas GMRES(120) does not converge in 2000 iterations. The savings in time is also noticeable. The gain ranges from two to four depending on the geometry and tends to become larger when the problem size increases. On the Cobra and the Almond test cases, the embedded solver reduces not only the solution time but also the memory used, as FGMRES(15)/GMRES(30) is faster than GMRES(120) (see Table 1).

Similarly to what has been observed with GMRES, using a full FGMRES in the outer loop enables us to reduce the number of outer iterations but often results in a longer solution time. In Table 7 we show results obtained with full FGMRES and compare this with the full GMRES method. It can also be seen that the savings in time due to the use of the inner-outer scheme with respect to GMRES is also strongly related to the reduction in the number of dot product calculations since there are far fewer basis vectors to be orthogonalized. This gain is particularly visible on the aircraft test example with 213084 unknowns. The time spent in the FMM is longer for the FGMRES/GMRES solver, about 3 hours and 47 minutes (i.e., $33 \times 11.8 + 1920 \times 6.9$; see Tables 7 and 4), and it is about 3 hours and 11 minutes for GMRES. In that example, the embedded scheme is about 2 hours faster than regular GMRES. In this latter case, we have to orthogonalize up to 973 basis vectors while in the inner-outer scheme we do not have to orthogonalize a basis larger than 60. This is another illustration that the orthogonalization procedure is extremely time consuming.

TABLE 6

Number of matrix-vector products and elapsed time required to converge on eight processors of the Compaq machine. The tests were run on eight processors of the Compaq machine, except those marked with ^(k), which were run on k processors.

Aircraft						
Size	Density FROB	Time FROB	GMRES(120)		FGMRES(30)/GMRES(60)	
			Iter	Time	Iter	Time
94704	0.28	11m	1956	3h 13m	27+1560	2h 14m
213084	0.13	31m	+2000	7h 56m	34+1920	5h
591900	0.09	1h 30m	+2000	1d 57m	57+3300	1d 9h 45m
1160124	0.02	3h 24m	+2000	> 4d	51+2940	16h 41m ⁽⁶⁴⁾
Almond						
Size	Density FROB	Time FROB	GMRES(60)		FGMRES(15)/GMRES(30)	
			Iter	Time	Iter	Time
104793	0.19	6m	302	19m	11+300	14m
419172	0.05	21m	+2000	> 10h	20+540	1h 24m
943137	0.02	49m	1633	6h 36m ⁽¹⁶⁾	22+600	3h 32m
Cetaf						
Size	Density FROB	Time FROB	GMRES(120)		FGMRES(30)/GMRES(60)	
			Iter	Time	Iter	Time
86256	0.18	4m	1546	1h 43m	17+ 960	55m
134775	0.11	6m	1125	1h 55m	15+ 840	1h 19m
264156	0.06	13m	1373	4h 46m	17+ 960	2h 22m
531900	0.03	20m	1717	14h 8m	19+1080	6h
1056636	0.01	37m	+2000	> 1d	22+1260	14h
Cobra						
Size	Density FROB	Time FROB	GMRES(60)		FGMRES(15)/GMRES(30)	
			Iter	Time	Iter	Time
60695	0.24	2m	708	29m	24+660	18m
179460	0.09	7m	433	48m	20+540	42m

TABLE 7

Number of matrix-vector products and elapsed time required to converge on eight processors of the Compaq machine. The tests were run on eight processors of the Compaq machine, except those marked with ^(k), that were run on k processors.

Aircraft				
Size	GMRES(∞)		FGMRES(∞ ,60)	
	Iter	Time	Iter	Time
94704	746	2h 9m	23+1320	2h 30m
213084	973	7h 19m	30+1740	6h 11m
591900	1461	16h 42m ⁽⁶⁴⁾	43+2520	12h ⁽³²⁾
1160124	M.L.E. ⁽⁶⁴⁾	> 1d	43+2520	14 h 28m ⁽⁶⁴⁾
Cobra				
Size	GMRES(∞)		FGMRES(∞ ,60)	
	Iter	Time	Iter	Time
60695	369	26m	21+600	17m
179460	353	1h 11m	18+510	38m

In Figure 3.3 we show the typical convergence history of the FGMRES/GMRES and GMRES solvers on a large Cetaf problem discretized with 264156 points. We depict the convergence curve as a function of the elapsed time. It can be seen that the embedded solver clearly outperforms the single GMRES scheme. The FGMRES/GMRES solver continues to succeed in reducing the residual norm while GMRES converges slowly. Similar behavior is also reported for simpler model examples in [29].

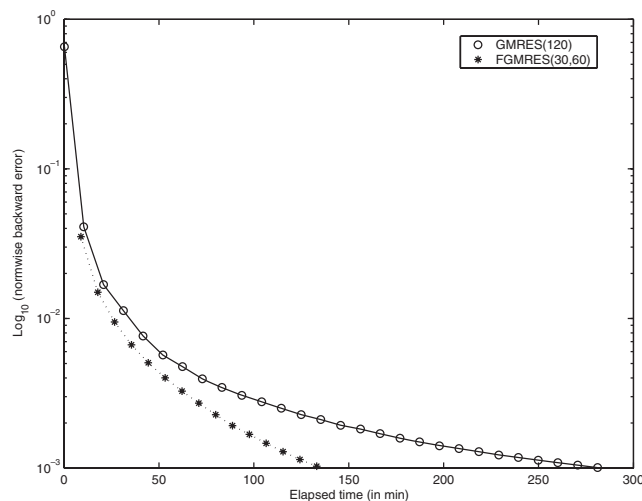


FIG. 3.3. Convergence history as a function of elapsed time of GMRES versus FGMRES on the Cetaf problem discretized with 264156 points.

For all the numerical experiments reported so far, we consider a discretization using ten points per wavelength. In that context, we vary the size of the linear system by varying the frequency of the illuminating wave and we consider objects with a fixed dimension expressed in meters. Consequently, each linear system corresponds to the simulation of a different physical phenomenon. In contrast, in Table 8 we solve the same physical phenomenon. We vary the size of the linear systems by varying the number of points per wavelength, keeping the frequency of the illuminating wave constant as well as the dimension of the object. In that table, λ denotes the wavelength and h the average mesh size. Because the leaf boxes used to build the preconditioner have a size that is a fraction of λ , the preconditioner becomes denser as the problem size increases. This observation partially explains the fact that the number of iterations of GMRES(60) slightly decreases when the size of the linear system increases. This is an opposite behavior to what we observed in Table 6, where the density of the preconditioner decreases when the size of the problem is increased. We also observe that the behavior of FGMRES(15)/GMRES(30) is not affected by the mesh refinement.

4. Conclusions. In this paper we have described an effective and inherently parallel approximate inverse preconditioner, based on a Frobenius-norm minimization preconditioner, that can easily be implemented in an FMM code. We have studied the numerical scalability of our method for the solution of very large and dense linear systems of equations arising from real-life applications in electromagnetism. The locality of the preconditioner on large problems can be significantly improved by embedding it within inner-outer solvers. We have described an embedded iterative scheme based on the GMRES method, implemented in a multipole context with different levels of accuracy for the matrix-vector products in the inner and outer loops. We have shown that the combination FGMRES/GMRES can effectively enhance the robustness of the preconditioner and reduce significantly the computational cost and the storage requirements for the solution of large problems. One could apply this idea recursively and embed several FGMRES schemes with decreasing FMM accuracy down to the

TABLE 8

Number of matrix-vector products and elapsed time required to converge on eight processors of the Compaq machine when the discretization is refined. The tests were run on eight processors of the Compaq machine, except those marked with $(^k)$, that were run on k processors.

Almond - length = 2.5 m - frequency 2.6 GHz					
h	Size of linear system	GMRES(60)		FGMRES(15,30)	
		Iter	Time	Iter	Time
$\lambda/6$	38724	345	15m	12+330	11m
$\lambda/10$	104793	302	19m	11+300	14m
$\lambda/15$	240294	281	1h 21m	10+270	1h 06m
$\lambda/20$	426750	271	2h 36m ⁽¹⁶⁾	10+270	2h 45m ⁽¹⁶⁾

lowest accuracy in the innermost GMRES. However, in our work, we consider only a two-level scheme. Most of the experiments shown in this paper require a huge amount of computation and storage, and they often reach the limits of our target machine in terms of memory and disk storage. To give an idea of how large these simulations are, the solution of systems with one million unknowns using a direct method would require 8 Tbytes of storage and 37 years of computation on one processor of the target computer (assuming the computation runs at peak performance). Such simulations are nowadays feasible thanks to the use of iterative methods and can be integrated in the design processes where the bottleneck moves from the simulation to the pre- and postprocessing of the results, as the tools are not yet available to easily manipulate meshes with millions of degrees of freedom. Finally we mention that the techniques described in this paper have recently been applied successfully to large simulations in computational acoustics and extended to the computation of monostatic radar cross sections [16].

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